

Quantum trajectories and non i.i.d. random products of matrices

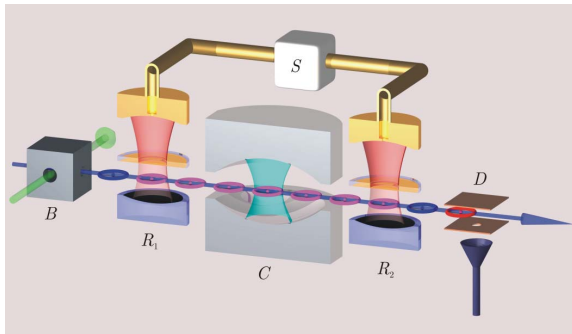
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A canonical experiment

S. Haroche group experiment:



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j 1101111111110011101101111  
i ddcbccabcdaadaabaddbadbc  
j 0101001101010101101011111  
i dababbaacbcccaddccdcbaaacc
```

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j 0001000110110000001010110  
i ddcaddabbccdccbcdaabbccab  
j 0001010100000100011101101  
i bcdaddaabbdbbdcdccadaada
```

Pictures: LKB ENS

Average evolution

On finite dimensional Hilbert spaces.

- ▶ **Quantum states:** Finite dimensional quantum system state: Density matrices \mathcal{D} :

$$\mathcal{D} = \{\rho \in M_d(\mathbb{C}) \mid \rho \geq 0, \quad \text{tr } \rho = 1\}.$$

- ▶ **Observables:** Self-adjoint matrices: $M_d^{sa}(\mathbb{C})$. To each observable A corresponds a random variable X_A on $\Omega = \{1, \dots, d\}$ whose moments are

$$\mathbb{E}_\rho(X_A^n) = \text{tr}(A^n \rho), \quad \forall n \in \mathbb{N}.$$

- ▶ **(Average) Evolution:** A completely positive trace preserving (CPTP) map $\Phi : \mathcal{D} \rightarrow \mathcal{D}$.

$$\exists \{V_i\}_{i=1, \dots, \ell} \subset M_d(\mathbb{C}), \quad \sum_{i=1, \dots, \ell} V_i^* V_i = I_d$$

such that,

$$\Phi(\rho) = \sum_{i=1}^{\ell} V_i \rho V_i^*.$$

$$\mathbb{E}(\rho_n) = \Phi^n(\rho).$$

Quantum trajectories

Definition (Unraveling)

Given a CPTP map Φ , fix one of its Kraus family $\{V_i\}_{i=1,\dots,\ell}$, then the stochastic process defined by

$$\rho_n = \frac{V_{i_n} \cdots V_{i_1} \rho V_{i_1}^* \cdots V_{i_n}^*}{\text{tr}(V_{i_n} \cdots V_{i_1} \rho V_{i_1}^* \cdots V_{i_n}^*)}, \quad \text{with proba. } \mathbb{P}_\rho(i_1, \dots, i_n) := \text{tr}(V_{i_n} \cdots V_{i_1} \rho V_{i_1}^* \cdots V_{i_n}^*)$$

is called an unraveling of Φ .

Definition (Quantum trajectory)

Let $\{V_i\}_{i=1,\dots,\ell} \subset M_d(\mathbb{C})$ define a CPTP map Φ . Then the Markov process defined by the Kernel,

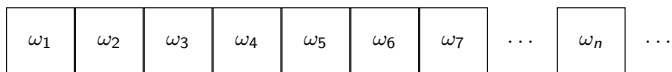
$$(\Pi f)(\rho) = \sum_{i=1}^{\ell} f\left(\frac{V_i \rho V_i^*}{\text{tr}(V_i \rho V_i^*)}\right) \text{tr}(V_i \rho V_i^*)$$

is called a Quantum trajectory.

Proposition

Quantum trajectories and Unravelings define the same processes and $\mathbb{E}_\rho(\rho_n) = \Phi^n(\rho)$.

The dynamical system picture



$\omega_n = 1, \dots, \ell$.

- ▶ States of the dynamical system: $\mathcal{A} := \{1, \dots, \ell\}$
- ▶ “Trajectory” space: $\Omega = \{1, \dots, \ell\}^{\mathbb{N}} \equiv [0, 1]$, time: $n \in \mathbb{N}^*$.
- ▶ Probability measure on the “trajectory” space Ω :

$$\mathbb{P}_\rho(\{\omega | \omega_k = i_k, 1 \leq k \leq n\}) = \text{tr}(V_{i_n} \cdots V_{i_1} \rho V_{i_1}^* \cdots V_{i_n}^*),$$

- ▶ Time shift: $f \circ \phi^n(\omega_1, \omega_2, \dots) = f(\omega_{n+1}, \omega_{n+2}, \dots)$.

$$\mathbb{P}_\rho \circ \phi^{-n} = \mathbb{P}_{\Phi^n(\rho)}.$$

Perron–Frobenius Theorem for CP maps

Theorem (Perron–Frobenius [Evans, Høegh-Krohn '77])

Let $\Phi : M_d(\mathbb{C}) \rightarrow M_d(\mathbb{C})$ be a CPTP map. Then the two following statements are equivalent.

- (i) **(Irr.)** If an orthogonal projector P is such that $V_i P \mathbb{C}^d \subset P \mathbb{C}^d$ for all $i \in \{1, \dots, \ell\}$, then $P \in \{0, I_d\}$.
- (ii) 1 is simple eigenvalue of Φ and the corresponding eigenstate is positive definite: $\exists! \rho_{inv.} \in \mathcal{D}$, s.t. $\Phi(\rho_{inv.}) = \rho_{inv.} > 0$.

Moreover, (i) or (ii) imply the peripheral spectrum of Φ is a finite subgroup of $U(1)$:

$$\text{spec}(\Phi) \cap U(1) = \{e^{i2\pi \frac{k}{m}}\}_{k=1, \dots, m}.$$

There also exist a unitary $U \in M_d(\mathbb{C})$ with spectral decomposition

$$U := \sum_{k=1}^m e^{i2\pi \frac{k}{m}} P_k$$

such that $\Phi(P_k) = P_{k-1}$.

Proposition (Convergence in total variation)

Assume Φ is irreducible. Then there exists $m \in \{1, \dots, d^2\}$, $C > 0$ and $\lambda < 1$ such that

$$\sup_{\rho \in \mathcal{D}} \sup_{A \subset \Omega} \left| \frac{1}{m} \sum_{k=1}^m \mathbb{P}_{\rho \circ \phi^{-k-mn}}(A) - \mathbb{P}_{\rho_{inv.}}(A) \right| \leq C\lambda^n.$$

Corollary

Assume Φ is irreducible and $m = 1$. Then the dynamical system $(\Omega, \phi, \mathbb{P}_{\rho_{inv.}})$ is exponentially mixing. For any measurable $A, B \subset \Omega$,

$$|\mathbb{P}_{\rho_{inv.}}(A\phi^{-n}(B)) - \mathbb{P}_{\rho_{inv.}}(A)\mathbb{P}_{\rho_{inv.}}(B)| \leq C\lambda^n.$$

Remark ([Guta, van Horsen '14; Carbone, Pautrat '15])

Using particularly Perron–Frobenius Theorem, the Law of Large Numbers, the Central Limit Theorem and a Large Deviation Principle follows directly for any random variable depending only on finite sequences of $\{1, \dots, \ell\}$ elements. A Large Deviation Principle also holds for the empirical measure over Ω .

Markov Chain first asymptotic properties

Theorem (First Law of Large Numbers [Kümmerer, Maassen '04])

Let $(\rho_n)_n$ be an unraveling of Φ . Then

$$\rho_\infty := \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=1}^n \rho_k$$

exists \mathbb{P}_ρ -almost surely and is such that $\Phi(\rho_\infty) = \rho_\infty$.

If moreover Φ is irreducible, $\rho_\infty = \rho_{\text{inv}}$. \mathbb{P}_ρ -almost surely.

Theorem (Purification [Kümmerer, Maassen '04])

Let $\{V_i\}_{i=1, \dots, \ell}$ be a finite family of $d \times d$ complex matrices corresponding to the Kraus decomposition of a CPTP map Φ .

(Pur.) Assume that any orthogonal projector Q such that $QV_i^*V_iQ \propto Q$ for all $i \in \{1, \dots, \ell\}$ is of rank 1.

Then the sequence $(\rho_n)_n$ purifies almost surely as $n \rightarrow \infty$. Namely, there almost surely exists a sequence of rank 1 orthogonal projectors $(|x_n\rangle\langle x_n|)_n \subset \mathcal{D}$ such that

$$\lim_{n \rightarrow \infty} \|\rho_n - |x_n\rangle\langle x_n|\| = 0, \quad \mathbb{P}_\rho - \text{a.s.}$$

It implies,

$$\lim_{n \rightarrow \infty} S(\rho_n) = \lim_{n \rightarrow \infty} -\text{tr}(\rho_n \ln \rho_n) = 0, \quad \mathbb{P}_\rho - \text{a.s.}$$

Non demolition measurement and wave function collapse

Theorem (Bauer, Bernard '11)

Non demolition: Assume that there exists an o.n.b. of pointers \mathcal{P} such that all the V_i 's are diagonal in \mathcal{P} .

Distinguishability: Assume that for any two different $x, y \in \mathcal{P}$, there exists i such that $\|V_i x\|_2 \neq \|V_i y\|_2$.

Then,

$$\lim_{n \rightarrow \infty} \rho_n = |\hat{x}\rangle\langle\hat{x}|, \text{ a.s.}$$

with $\hat{x} : \Omega \rightarrow \mathcal{P}$ and

$$\mathbb{P}_\rho(\hat{x} = y) = \text{tr}(|y\rangle\langle y|\rho).$$

Ergodic properties of the Markov Chain

Theorem (B., Fraas, Pautrat, Pellegrini '16)

Assume **(Irr.)** and **(Pur.)** hold. Then, Π accepts a unique invariant measure $\nu_{inv.}$ and there exists $m \in \{1, \dots, d^2\}$, $C > 0$ and $\lambda < 1$ such that, for any measure ν over \mathcal{D} ,

$$W_1 \left(\frac{1}{m} \sum_{k=1}^m \nu \Pi^{k+mn}, \nu_{inv.} \right) \leq C \lambda^n$$

where W_1 is the order 1 Wasserstein metric.

Remark

- ▶ No convergence in total variation possible since if ν_a has a continuous support, ν_b is pure point and all the V_i 's are invertible,

$$\|\nu_a \Pi^n - \nu_b \Pi^n\|_{TV} = 1, \quad \forall n \in \mathbb{N}.$$

- ▶ ϕ -irreducibility methods are not suitable for this Markov Chain.
- ▶ Proof inspired by product of i.i.d. random matrices.
- ▶ Under stronger conditions (invertibility, strong irreducibility) Guivarc'h and Le Page proved a similar result in 2004.
- ▶ Assumptions **(Irr.)** and **(Pur.)** are optimal. A slightly more general assumption **(Irr.)** is necessary.

Proof structure

Proof.

1. Define $\hat{\mu}_n$ as a sequence of estimates of the initial state given the growing sequence $(\omega_1, \dots, \omega_n)$:

$$\hat{\mu}_n = \operatorname{argmax}_{\mu \in \mathcal{D}} \ln \mathbb{P}_\mu(\omega_1, \dots, \omega_n).$$

Let

$$\hat{\rho}_n := \frac{V_{i_n} \cdots V_{i_1} \hat{\mu}_n V_{i_1}^* \cdots V_{i_n}^*}{\operatorname{tr}(V_{i_n} \cdots V_{i_1} \hat{\mu}_n V_{i_1}^* \cdots V_{i_n}^*)}.$$

Then **(Pur.)** implies

$$\mathbb{E}_\nu(d(\hat{\rho}_n, \rho_n)) \leq C\lambda^n.$$

2. The estimated sequence $(\hat{\rho}_n)$ not depending on the initial state ρ , the convergence follows from the ergodic properties of $\mathbb{P}_{\mathbb{E}_\nu(\rho)}$ implied by **(Irr.)**, the uniformity of the sub geometric bounds in the initial state ρ and the Markov property.
3. The Wasserstein metric bound is obtained by Kantorovich and Rubinstein Duality Theorem.

□

First consequences

Theorem (Law of Large Numbers [BFPP '16])

Assume **(Irr.)** and **(Pur.)** hold. Then for any continuous function on \mathcal{D} , for any initial probability measure ν over \mathcal{D} ,

$$\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=1}^n f(\rho_k) = \mathbb{E}_{\nu_{inv.}}(f(\rho)), \quad \text{a.s.}$$

Theorem (Functional Central Limit Theorem [BFPP '16])

Assume **(Irr.)** and **(Pur.)** hold. Then for any Hölder continuous function g the Functional Central Limit Theorem Holds.

Current developments

- ▶ **Hypothesis testing** Ability to distinguish between different CPTP maps (mutual singularity and error exponents).
Applied to Hypothesis Testing of the arrow of time [B., Jaksic, Pautrat, Pillet '16].
- ▶ **Parameter estimation** [Guta, Kiukas, Levitt '15–'16].
- ▶ **Dynamical Phases** Characterisation of the existence of dynamical phases[Guta, van Horssen '14].
Full characterisation in terms of non differentiability of Rényi entropy on \mathbb{R}_+ [BJPP '16 in preparation].
Link with selection of invariant states and the failure of the CLT[B., Pautrat, Pellegrini in preparation].
First approach to the characterisation of metastable behaviour [Macieszczak, Guta, Lesanovsky, Garrahan '16].

Markov Chain

- ▶ Regularity of the invariant measure ν_{inv} .
- ▶ Relaxation of **(Pur.)**. Mixing for products of i.i.d. elements of $SU(d)$?
- ▶ Large Deviation Principle for the chain. Existence of a spectral gap for Π .
- ▶ Entropy production of the chain and time reversibility.

Dynamical system

- ▶ Meaning of the irregularities of the Rényi entropy in terms of dynamical phases beyond non differentiability on \mathbb{R}_+ .
- ▶ Full characterisation of Dynamical Phase Transitions and metastability.

A problem for continuous time quantum trajectories

In the appropriate continuous time limit, discrete quantum trajectories converge weakly towards solutions of stochastic differential equations[Pellegrini '08-'10].

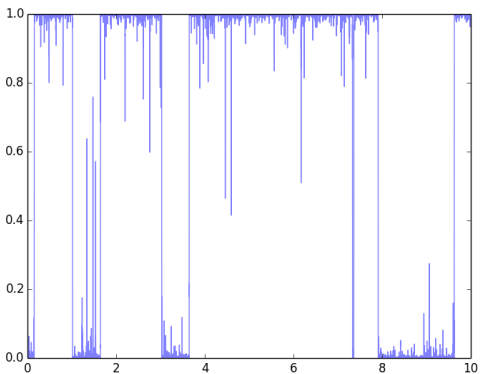
One particular issue arise in the emergence of quantum jumps[Bauer,Bernard,Tilloy '15].

The energy population p_t of the excited state of a two level atom being indirectly measured by a diffusive signal is a process solution of the following diffusive SDE.

$$dp_t = \lambda\left(\frac{1}{2} - p_t\right)dt + \sqrt{\gamma}p_t(1 - p_t)dW_t, \quad p_0 \in (0, 1).$$

What is the limit of $(p_t)_t$ when $\gamma \rightarrow \infty$?

Numerical simulation of $(p_t)_t$ ($\lambda = 1, \gamma = 10^4$):



The “spikes” do not disappear when the simulation is refined.

How to make sense of the limit?

- ▶ Use an appropriate topology that does not see the “spikes” (Meyer–Zheng convergence in measure topology).
- ▶ Find a meaning full limit towards a random variable taking value in nowhere continuous functions.
- ▶ Limit towards a Poisson measure on $[0, 1] \times [0, T]$ for the local maxima.

Another approach [Bauer, Bernard, Tilloy '16]:

Study in effective time by Dubins-Schwarz Theorem. Let

$$\tau_s := \inf \left\{ t : \gamma \int_0^t p_u^2 (1 - p_u)^2 du > s \right\}.$$

Then $\tilde{p}_s := p_{\tau_s}$ should be the solution of

$$d\tilde{p}_s = \frac{\frac{1}{2} - \tilde{p}_s}{\gamma \tilde{p}_s^2 (1 - \tilde{p}_s)^2} ds + dB_s, \quad \tilde{p}_0 = p_0.$$

[Bauer, Bernard, Tilloy '16] provides a formal proof that $(\tilde{p}_s)_s$ converges towards a Brownian motion reflected in 0 and 1.

Thank you!